



ELSEVIER

Available online at www.sciencedirect.com
 Communications in
 Nonlinear Science and
 Numerical Simulation

Communications in Nonlinear Science and Numerical Simulation xxx (2006) xxx–xxx

www.elsevier.com/locate/cnsns

Invariants and invariant description of second-order ODEs with three infinitesimal symmetries. II

 Nail H. Ibragimov ^a, Sergey V. Meleshko ^{b,*}
^a *Department of Mathematics and Science Research Centre ALGA: Advances in Lie Group Analysis, Blekinge Institute of Technology, SE-371 79 Karlskrona, Sweden*
^b *School of Mathematics, Suranaree University of Technology, Nakhon Ratchasima 3000, Thailand*

Received 18 February 2006; accepted 18 March 2006

Abstract

The second-order ordinary differential equations can have one, two, three or eight independent symmetries. Sophus Lie showed that the equations with eight symmetries and only these equations can be linearized by a change of variables. Moreover he demonstrated that these equations are at most cubic in the first derivative and gave a convenient invariant description of all linearizable equations. We provide a similar description of the equations with three symmetries. There are four different types of such equations. Classes of equations equivalent to one of these equations were studied in [Ibragimov NH, Meleshko SV. Invariants and invariant description of second-order ODEs with three infinitesimal symmetries. Communication in Nonlinear Science and Numerical Simulation, in press], where we presented the candidates for all four types and studied one of these candidates. The present paper is the continuation of the work of Ibragimov and Meleshko and is devoted to other three candidates.

© 2006 Elsevier B.V. All rights reserved.

PACS: 02.30.Jr; 02.20.Tw

Keywords: Equations with three symmetries; Candidates; Equivalence test

1. Introduction

According to Lie's classification [2] in the complex domain, any ordinary differential equation of the second order

$$y'' = f(x, y, y'), \quad (1)$$

admitting a three-dimensional Lie algebra belongs to one of four distinctly different types. Each of these four types is obtained by a change of variables from the following canonical representatives (see, e.g. [3], Section 8.4):

* Corresponding author. Tel.: +66 44 224382; fax: +66 44 224185.
 E-mail address: sergey@math.sut.ac.th (S.V. Meleshko).

$$y'' + Cy^{-3} = 0, \quad (2)$$

$$y'' + e^{y'} = 0, \quad (3)$$

$$y'' + y^{(k-2)/(k-1)} = 0, \quad (4)$$

$$y'' + 2 \frac{y' + Cy^{3/2} + y'^2}{x - y} = 0, \quad (5)$$

where k and C are constants such that $k \neq 0, 1/2, 1, 2$ and $C \neq 0$.

Eqs. (2)–(5) admit non-similar three-dimensional Lie algebras L_3 spanned by the operators

$$X_1 = \frac{\partial}{\partial x}, \quad X_2 = 2x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y}, \quad X_3 = x^2 \frac{\partial}{\partial x} + xy \frac{\partial}{\partial y}, \quad (6)$$

$$X_1 = \frac{\partial}{\partial x}, \quad X_2 = \frac{\partial}{\partial y}, \quad X_3 = x \frac{\partial}{\partial x} + (y - x) \frac{\partial}{\partial y}, \quad (7)$$

$$X_1 = \frac{\partial}{\partial x}, \quad X_2 = \frac{\partial}{\partial y}, \quad X_3 = x \frac{\partial}{\partial x} + ky \frac{\partial}{\partial y}, \quad (8)$$

and

$$X_1 = \frac{\partial}{\partial x} + \frac{\partial}{\partial y}, \quad X_2 = x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y}, \quad X_3 = x^2 \frac{\partial}{\partial x} + y^2 \frac{\partial}{\partial y}, \quad (9)$$

respectively (see, e.g. [3], Section 8.4).

2. Candidates for equations with three symmetries

Let us subject each of Eqs. (2)–(5) to the arbitrary change of variables

$$t = \varphi(x, y), \quad u = \psi(x, y), \quad (10)$$

where t is a new independent variable and u is a new dependent variable. Then we obtain from (2)–(5) the equations of the form

$$u'' + b_1 u'^3 + 3b_2 u'^2 + 3b_3 u' + b_4 = 0, \quad (11)$$

$$u'' + b_1 u'^3 + 3b_2 u'^2 + 3b_3 u' + b_4 + (b_5 u'^3 + 3b_6 u'^2 + 3b_7 u' + b_8) \exp\left(\frac{b_9 u' + b_{10}}{b_{11} u' + b_{12}}\right) = 0, \quad (12)$$

$$u'' + b_1 u'^3 + 3b_2 u'^2 + 3b_3 u' + b_4 + (b_5 u'^3 + 3b_6 u'^2 + 3b_7 u' + b_8) \left(\frac{b_9 u' + b_{10}}{b_{11} u' + b_{12}}\right)^{(k-2)/(k-1)} = 0, \quad (13)$$

and

$$u'' + b_1 u'^3 + 3b_2 u'^2 + 3b_3 u' + b_4 + (b_5 u'^3 + 3b_6 u'^2 + 3b_7 u' + b_8) C \left(\frac{b_9 u' + b_{10}}{b_{11} u' + b_{12}}\right)^{3/2} = 0, \quad (14)$$

respectively, where $b_i = b_i(t, u)$, $i = 1, \dots, 12$. Eqs. (11)–(14) are the candidates for the equations with three symmetries.

All candidates can be encapsulated in the formula

$$u'' + b_1 u'^3 + 3b_2 u'^2 + 3b_3 u' + b_4 + (b_5 u'^3 + 3b_6 u'^2 + 3b_7 u' + b_8) f\left(\frac{b_9 u' + b_{10}}{b_{11} u' + b_{12}}\right) = 0.$$

Namely, Eqs. (11)–(14) are obtained by letting

$$f(z) = 0, \quad f(z) = e^z, \quad f(z) = z^{(k-2)/(k-1)}, \quad f(z) = z^{3/2}. \quad (15)$$

Using the usual formula for the transformation of derivatives under the change of variables (10), we obtain the following statement.

Theorem. Any equation of the form

$$u'' + b_1u'^3 + 3b_2u'^2 + 3b_3u' + b_4 + (b_5u'^3 + 3b_6u'^2 + 3b_7u' + b_8)f\left(\frac{b_9u' + b_{10}}{b_{11}u' + b_{12}}\right) = 0 \tag{16}$$

is transformed by the change of variables (10) into an equation of the same form:

$$y'' + a_1y'^3 + 3a_2y'^2 + 3a_3y' + a_4 + (a_5y'^3 + 3a_6y'^2 + 3a_7y' + a_8)f\left(\frac{a_9y' + a_{10}}{a_{11}y' + a_{12}}\right) = 0. \tag{17}$$

Here a_i and b_i are functions of x, y and t, u , respectively, and are connected by:

$$\begin{aligned} a_1 &= \Delta^{-1}[\varphi_y\psi_{yy} - \varphi_{yy}\psi_y + b_4\varphi_y^3 + 3b_3\varphi_y^2\psi_y + 3b_2\varphi_y\psi_y^2 + b_1\psi_y^3], \\ a_2 &= \Delta^{-1}[b_4\varphi_x\varphi_y^2 + b_3\varphi_x(2\varphi_x\psi_y + \varphi_y\psi_x) + b_2\psi_y(\varphi_x\psi_y + 2\varphi_y\psi_x) \\ &\quad + b_1\psi_x\psi_y^2 + (\varphi_x\psi_{yy} - \varphi_{yy}\psi_x - 2\varphi_{xy}\psi_y + 2\varphi_y\psi_{xy})/3], \end{aligned} \tag{18}$$

$$\begin{aligned} a_3 &= \Delta^{-1}[b_4\varphi_x^2\varphi_y + b_3\varphi_x(\varphi_x\psi_y + 2\varphi_y\psi_x) + b_2\psi_x(2\varphi_x\psi_y + \varphi_y\psi_x) \\ &\quad + b_1\psi_x^2\psi_y + (\varphi_y\psi_{xx} - \varphi_{xx}\psi_y - 2\varphi_{xy}\psi_x + 2\varphi_x\psi_{xy})/3], \end{aligned} \tag{19}$$

$$\begin{aligned} a_4 &= \Delta^{-1}[b_4\varphi_x^3 + 3b_3\varphi_x^2\psi_x + 3b_2\varphi_x\psi_x^2 + b_1\psi_x^3 - \varphi_{xx}\psi_x + \varphi_x\psi_{xx}], \\ a_5 &= \Delta^{-1}[b_8\varphi_y^3 + 3b_7\varphi_y^2\psi_y + 3b_6\varphi_y\psi_y^2 + b_5\psi_y^3], \\ a_6 &= \Delta^{-1}[b_8\varphi_x\varphi_y^2 + b_7\varphi_x(2\varphi_x\psi_y + \varphi_y\psi_x) + b_6\psi_y(\varphi_x\psi_y + 2\varphi_y\psi_x) + b_5\psi_x\psi_y^2], \end{aligned} \tag{19}$$

$$\begin{aligned} a_7 &= \Delta^{-1}[b_8\varphi_x^2\varphi_y + b_7\varphi_x(\varphi_x\psi_y + 2\varphi_y\psi_x) + b_6\psi_x(2\varphi_x\psi_y + \varphi_y\psi_x) + b_5\psi_x^2\psi_y], \\ a_8 &= \Delta^{-1}[b_8\varphi_x^3 + 3b_7\varphi_x^2\psi_x + 3b_6\varphi_x\psi_x^2 + b_5\psi_x^3], \end{aligned} \tag{20}$$

$$a_9 = b_{10}\varphi_y + b_9\psi_y, \quad a_{10} = b_{10}\varphi_x + b_9\psi_x, \quad a_{11} = b_{12}\varphi_y + b_{11}\psi_y, \quad a_{12} = b_{12}\varphi_x + b_{11}\psi_x, \tag{20}$$

where

$$\Delta = (\varphi_x\psi_y - \varphi_y\psi_x) \neq 0$$

is the Jacobian of the change of variables (10).

It follows from Eqs. (20) that

$$a_9a_{12} - a_{10}a_{11} = \Delta(b_9b_{12} - b_{10}b_{11}).$$

Hence the equation

$$a_9a_{12} - a_{10}a_{11} = 0 \tag{21}$$

is invariant under the change of variables (10). If $a_9a_{12} - a_{10}a_{11} = 0$, and hence $b_9b_{12} - b_{10}b_{11} = 0$, the function f disappears in both Eqs. (16) and (17). This leads to the equations equivalent to Eq. (2), i.e. to the case considered in [1]. Therefore, we assume in what follows that

$$b_9b_{12} - b_{10}b_{11} \neq 0, \quad a_9a_{12} - a_{10}a_{11} \neq 0.$$

3. Equations equivalent to Eqs. (3) and (4)

The test for equivalence to both Eqs. (3) and (4) have the same form. The only difference is that Eqs. (3) and (4) have the candidates (12) and (13), respectively, with the different functions f .

Eqs. (3) and (4) have the form (17) with

$$\begin{aligned} a_1 &= 0, & a_2 &= 0, & a_3 &= 0, & a_4 &= 0, & a_5 &= 0, & a_6 &= 0, \\ a_7 &= 0, & a_8 - 1 &= 0, & a_9 - a_{12} &= 0, & a_{10} &= 0, & a_{11} &= 0, \end{aligned} \tag{22}$$

whereas the function f has the form $f(z) = e^z$ for (3) and $f(z) = z^{(k-2)/(k-1)}$ for (4). Furthermore, the change of variables (10) leaves invariant each candidate. Hence, the equations which are equivalent to (3) and (4) belong to equations of the form (16):

$$u'' + b_1u'^3 + 3b_2u'^2 + 3b_3u' + b_4 + (b_5u'^3 + 3b_6u'^2 + 3b_7u' + b_8)f\left(\frac{b_9u' + b_{10}}{b_{11}u' + b_{12}}\right) = 0.$$

Thus for the functions (10) $\varphi(x, y)$ and $\psi(x, y)$ one obtains the overdetermined system of equations which consists of Eqs. (22), where the coefficients a_i , ($i = 1, 2, \dots, 12$) are defined by the relations (18)–(20).

Analysis of compatibility of the overdetermined system depends on the value of b_{12} . If the argument of the function f in (16) is a linear function with respect to the derivative u' , then without loss of generality one can assume that $b_{11} = 0$ and $b_{12} = 1$. If the argument of the function f in (16) is a rational function with respect to the derivative u' , then without loss of generality one can assume that $b_{11} = 1$.

Let us consider the first case

$$b_{11} = 1, \quad b_{12} = 0.$$

In this case the result of compatibility analysis gives that $b_5 b_{10} \neq 0$ and

$$\begin{aligned} b_4 &= 0, & b_6 &= 0, & b_7 &= 0, & b_8 &= 0, \\ b_{5t} - 3b_5 b_3 &= 0, & b_{10} b_{5u} - 3b_5(2b_{10} b_2 + b_{9t}) &= 0, & b_{9u} + b_{10} b_1 &= 0, \\ b_{10t} + 3b_{10} b_3 &= 0, & b_{10u} + 3b_{10} b_2 + b_{9t} &= 0. \end{aligned}$$

The functions $\varphi(x, y)$ and $\psi(x, y)$ are found from the compatible system of equations

$$\varphi_x = \frac{b_9}{b_{10}^2 b_5}, \quad \varphi_y = -\frac{1}{b_{10}^2 b_5}, \quad \psi_x = -\frac{1}{b_{10} b_5}, \quad \psi_y = 0.$$

The generators corresponding to (7) are

$$\begin{aligned} X_1 &= (b_{10}^2 b_5)^{-1} \left[b_9 \frac{\partial}{\partial t} - b_{10} \frac{\partial}{\partial u} \right], & X_2 &= (b_{10}^2 b_5)^{-1} \frac{\partial}{\partial t}, \\ X_3 &= (b_{10}^2 b_5)^{-1} \left[((b_9 + 1)x - y) \frac{\partial}{\partial t} - b_{10} x \frac{\partial}{\partial u} \right]. \end{aligned}$$

The generators corresponding to (8) are

$$\begin{aligned} X_1 &= (b_{10}^2 b_5)^{-1} \left[b_9 \frac{\partial}{\partial t} - b_{10} \frac{\partial}{\partial u} \right], & X_2 &= (b_{10}^2 b_5)^{-1} \frac{\partial}{\partial t}, \\ X_3 &= (b_{10}^2 b_5)^{-1} \left[(b_9 x - ky) \frac{\partial}{\partial t} - b_{10} x \frac{\partial}{\partial u} \right]. \end{aligned}$$

In the second case $b_{12} = 1$ one obtains that $b_8 \neq 0$ and

$$\begin{aligned} b_{11t} b_{11} - b_{11u} + b_{11}^3 b_4 - 3b_{11}^2 b_3 + 3b_{11} b_2 - b_1 &= 0, \\ 3b_2(b_{11} b_{10} - b_9) b_8 + (b_{11} b_{10} - b_9) b_{8u} + 3b_{11} b_8 b_{10u} &= 0, \\ 6b_3(b_{11} b_{10} - b_9) b_8 + (b_{11} b_{10} - b_9) b_{8t} + 3b_{11} b_8 b_{10t} + 3b_8 b_{10u} &= 0, \\ b_4(b_{11} b_{10} - b_9) + b_{10t} &= 0, & b_5 &= b_{11}^3 b_8, & b_6 &= b_{11}^2 b_8, & b_7 &= b_{11} b_8, \\ b_8(2b_{11t} b_{10} + b_{10t} b_{11} + b_{10u} - 2b_{9t}) + b_{8t}(b_{11} b_{10} - b_9) &= 0, \\ (b_{11} b_{10} - b_9)(2b_{8u} - 2b_{11t} b_8 - b_{8t} b_{11}) + b_8(2b_{11u} b_{10} - b_{10t} b_{11}^2 + 3b_{10u} b_{11} - 2b_{9u}) &= 0. \end{aligned}$$

The functions $\varphi(x, y)$ and $\psi(x, y)$ are found from the compatible system of equations

$$\begin{aligned} \varphi_x &= \frac{b_9}{b_8(b_{11} b_{10} - b_9)^2}, & \varphi_y &= -\frac{b_{11}}{b_8(b_{11} b_{10} - b_9)^2}, \\ \psi_x &= -\frac{b_{10}}{b_8(b_{11} b_{10} - b_9)^2}, & \psi_y &= \frac{1}{b_8(b_{11} b_{10} - b_9)^2}. \end{aligned}$$

The generators corresponding to (7) are

$$\begin{aligned} X_1 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[b_9 \frac{\partial}{\partial t} - b_{10} \frac{\partial}{\partial u} \right], \\ X_2 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[b_{11} \frac{\partial}{\partial t} - \frac{\partial}{\partial u} \right], \\ X_3 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[(b_{11}(x-y) + b_9x) \frac{\partial}{\partial t} + (y - (1 + b_{10})x) \frac{\partial}{\partial u} \right]. \end{aligned}$$

The generators corresponding to (8) are

$$\begin{aligned} X_1 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[b_9 \frac{\partial}{\partial t} - b_{10} \frac{\partial}{\partial u} \right], \\ X_2 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[b_{11} \frac{\partial}{\partial t} - \frac{\partial}{\partial u} \right], \\ X_3 &= b_8^{-1}(b_{11}b_{10} - b_9)^{-2} \left[(b_9x - b_{11}ky) \frac{\partial}{\partial t} + (ky - b_{10}x) \frac{\partial}{\partial u} \right]. \end{aligned}$$

4. Equations equivalent to Eq. (5)

Eq. (5) has the form (17) with

$$\begin{aligned} a_1 &= 0, & 3a_2 - 2/(x-y) &= 0, & 3a_3 - 2/(x-y) &= 0, \\ a_4 &= 0, & a_5 &= 0, & a_6 &= 0, & a_7 &= 0, & a_8 - C(x-y)^{-1} &= 0, \\ a_9 - a_{12} &= 0, & a_{10} &= 0, & a_{11} &= 0, \end{aligned} \quad (23)$$

and f has the form $f(z) = Cz^{3/2}$. The equations that are equivalent to (5) belong to equations of the form (16):

$$u'' + b_1u'^3 + 3b_2u'^2 + 3b_3u' + b_4 + \left(b_5u'^3 + 3b_6u'^2 + 3b_7u' + b_8 \right) f \left(\frac{b_9u' + b_{10}}{b_{11}u' + b_{12}} \right) = 0.$$

Thus for the functions (10) $\varphi(x,y)$ and $\psi(x,y)$ one obtains the overdetermined system of equations which consists of Eqs. (23), where the coefficients a_i ($i = 1, 2, \dots, 12$) are defined by the relations (18)–(20).

Analysis of compatibility of the overdetermined system depends on the value of b_{12} . If the argument of the function f in (16) is a linear function with respect to the derivative u' , then without loss of generality one can assume that $b_{11} = 0$ and $b_{12} = 1$. If the argument of the function f in (16) is a rational function with respect to the derivative u' , then without loss of generality one can assume that $b_{11} = 1$.

Let us consider the first case

$$b_{11} = 1, \quad b_{12} = 0.$$

In this case the result of compatibility analysis gives that $b_5b_{10} \neq 0$ and

$$\begin{aligned} b_4 &= 0, & b_6 &= 0, & b_7 &= 0, & b_8 &= 0, \\ b_{10}C - b_{10}(2b_{10}^2b_5 - 3b_3C) &= 0, \\ 4b_{10}^2b_9b_5 + 2b_{10}^2b_5 - 3b_{10}b_2C - b_9C - b_{10u}C &= 0, \\ 3b_5(-b_{10}^2b_5 + b_3C) - b_{5t}C &= 0, \\ 3b_5(-3b_{10}^2b_9b_5 - b_{10}^2b_5 + 2b_{10}b_2C + b_9C) - b_{5u}b_{10}C &= 0, \\ b_{10}(2b_9^2b_5 + 2b_9b_5 - b_1C) - b_{9u}C &= 0. \end{aligned}$$

The functions $\varphi(x, y)$ and $\psi(x, y)$ are found from the compatible system of equations

$$\varphi_x = \frac{b_9 C}{b_{10}^2 b_5 (x-y)}, \quad \varphi_y = -\frac{C}{b_{10}^2 b_5 (x-y)}, \quad \psi_x = -\frac{C}{b_{10} b_5 (x-y)}, \quad \psi_y = 0.$$

The generators corresponding to (9) are

$$\begin{aligned} X_1 &= (b_{10}^2 b_5 (x-y))^{-1} \left[(b_9 - 1) \frac{\partial}{\partial t} - b_{10} \frac{\partial}{\partial u} \right], \\ X_2 &= (b_{10}^2 b_5 (x-y))^{-1} \left[(b_9 x - y) \frac{\partial}{\partial t} - b_{10} x \frac{\partial}{\partial u} \right], \\ X_3 &= (b_{10}^2 b_5 (x-y))^{-1} \left[(b_9 x^2 - y^2) \frac{\partial}{\partial t} - b_{10} x^2 \frac{\partial}{\partial u} \right]. \end{aligned}$$

In the second case $b_{12} = 1$ one obtains that $b_8 \neq 0$ and

$$\begin{aligned} b_5 &= b_{11}^3 b_8, \quad b_6 = b_{11}^2 b_8, \quad b_7 = b_{11} b_8, \\ -b_{10r} C + (b_{11} b_{10} - b_9)(2b_{10} b_8 (b_{10} + 1) - b_4 C) &= 0, \\ b_{11t} b_{11} - b_{11u} + b_{11}^3 b_4 - 3b_{11}^2 b_3 + 3b_{11} b_2 - b_1 &= 0, \\ -3b_{10u} b_{11} b_8 C - b_{8u} C (b_{11} b_{10} - b_9) + 3b_8^2 (b_{11} + b_9)(b_{11}^2 b_{10}^2 - b_9^2) - 3(b_{11} b_{10} - b_9) b_8 b_2 C &= 0, \\ -3b_{10r} b_{11} b_8 C - 3b_{10u} b_8 C + b_{8r} C (-b_{11} b_{10} + b_9) + 3b_8^2 (b_{11}^2 b_{10}^3 + 3b_{11}^2 b_{10}^2 + 2b_{11} b_{10}^2 b_9 - 2b_{11} b_{10} b_9 \\ - 3b_{10} b_9^2 - b_9^2) - 6(b_{11} b_{10} - b_9) b_3 b_8 C &= 0, \\ 2b_{11t} b_{10} b_8 C + b_{10r} b_{11} b_8 C + b_{10u} b_8 C - 2b_{9r} b_8 C + b_{8r} C (b_{11} b_{10} - b_9) \\ + b_8^2 (b_{11}^2 b_{10}^3 - b_{11}^2 b_{10}^2 - 2b_{11} b_{10}^2 b_9 + 2b_{11} b_{10} b_9 + b_{10} b_9^2 - b_9^2) &= 0, \\ 2b_{11t} b_8 C (-b_{11} b_{10} + b_9) + 2b_{11u} b_{10} b_8 C - b_{10r} b_{11}^2 b_8 C + 3b_{10u} b_{11} b_8 C + b_{8r} b_{11} C (-b_{11} b_{10} + b_9) \\ + 2b_{8u} C (b_{11} b_{10} - b_9) + b_8^2 (-b_{11}^3 b_{10}^3 - b_{11}^3 b_{10}^2 + 4b_{11}^2 b_{10}^2 b_9 + 2b_{11}^2 b_{10} b_9 - 5b_{11} b_{10} b_9^2 \\ - b_{11} b_9^2 + 2b_9^3) - 2b_{9u} b_8 C &= 0 \end{aligned}$$

The functions $\varphi(x, y)$ and $\psi(x, y)$ are found from the compatible system of equations

$$\begin{aligned} \varphi_x &= \frac{b_9 C}{(b_{11} b_{10} - b_9)^2 (x-y) b_8}, \quad \varphi_y = -\frac{b_{11} C}{(b_{11} b_{10} - b_9)^2 (x-y) b_8}, \\ \psi_x &= -\frac{b_{10} C}{(b_{11} b_{10} - b_9)^2 (x-y) b_8}, \quad \psi_y = \frac{C}{(b_{11} b_{10} - b_9)^2 (x-y) b_8}. \end{aligned}$$

The generators corresponding to (9) are

$$\begin{aligned} X_1 &= \left((b_{11} b_{10} - b_9)^2 (x-y) b_8 \right)^{-1} \left[(b_9 - b_{11}) \frac{\partial}{\partial t} + (1 - b_{10}) \frac{\partial}{\partial u} \right], \\ X_2 &= \left((b_{11} b_{10} - b_9)^2 (x-y) b_8 \right)^{-1} \left[(b_9 x - b_{11} y) \frac{\partial}{\partial t} + (y - b_{10} x) \frac{\partial}{\partial u} \right], \\ X_3 &= \left((b_{11} b_{10} - b_9)^2 (x-y) b_8 \right)^{-1} \left[(b_9 x^2 - b_{11} y^2) \frac{\partial}{\partial t} + (y^2 - b_{10} x^2) \frac{\partial}{\partial u} \right]. \end{aligned}$$

References

- [1] Ibragimov NH, Meleshko SV. Invariants and invariant description of second-order ODEs with three infinitesimal symmetries. I. Communication in Nonlinear Science and Numerical Simulation, in press, doi:10.1016/j.cnsns.2005.12.012.
- [2] Lie S. Klassifikation und Integration von gewöhnlichen Differentialgleichungen zwischen x, y , die eine Gruppe von Transformationen gestatten. III. Archiv for Matematik og Naturvidenskab., 8, Heft 1883;4:371–458. Reprinted in Lie's Ges. Abhandl., vol. 5, 1924, paper XIV, p. 362–427.
- [3] Ibragimov NH, editor CRC Handbook of Lie group analysis of differential equations. Symmetries, exact solutions and conservation laws, vol. 1. Boca Roton: CRC Press Inc.; 1994.